
ON THE GRAY-HERVELLA CLASSES OF *AH*-STRUCTURES
ON SIX-DIMENSIONAL SUBMANIFOLDS
OF CAYLEY ALGEBRA

M. BANARU

It is proved that only eight Gray-Hervella classes of almost Hermitian structures can be represented on six-dimensional general-type submanifolds of the octave algebra.

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1. PRELIMINARIES

In this paper we consider the Gray-Hervella classes of almost Hermitian (*AH*) structures [11]. As it is well-known [1, 11], a $2n$ -dimensional manifold M^{2n} with a Riemannian metric $g = \langle \cdot, \cdot \rangle$ and an almost complex structure J is called almost Hermitian if the following condition holds:

$$\langle JX, JY \rangle = \langle X, Y \rangle, \quad \forall X, Y \in \mathfrak{N}(M^{2n}),$$

where $\mathfrak{N}(M^{2n})$ is the module of smooth vector fields on M^{2n} . All considered manifolds, tensor fields and similar objects are assumed to be of the class C^∞ . We recall that the fundamental form of an almost Hermitian manifold is determined by

$$F(X, Y) = \langle X, JY \rangle, \quad X, Y \in \mathfrak{N}(M^{2n}).$$

The specification of an almost Hermitian structure on a manifold is equivalent to the setting of a G -structure, where G is the unitary group $U(n)$ [1]. Its elements

are the frames adapted to the structure (A -frames). They look as follows:

$$(p, \varepsilon_1, \dots, \varepsilon_n, \varepsilon_{\hat{1}}, \dots, \varepsilon_{\hat{n}}),$$

where $p \in M^{2n}$, ε_a are the eigenvectors corresponding to the eigenvalue $i = \sqrt{-1}$, and $\varepsilon_{\hat{a}}$ are the eigenvectors corresponding to the eigenvalue $-i$. Here $a = 1, \dots, n$; $\hat{a} = a + n$. Therefore, the matrix of the almost complex structure in an A -frame at the point p looks as follows:

$$(J_j^k) = \left(\begin{array}{c|c} iI_n & 0 \\ \hline 0 & -iI_n \end{array} \right), \quad (1.1)$$

where I_n is the identity matrix; $k, j = 1, \dots, 2n$. By direct computing, it is easy to obtain that in A -frame the matrices of the Riemannian metric g and of the fundamental form F look as follows, respectively:

$$(g_{kj}) = \left(\begin{array}{c|c} 0 & I_n \\ \hline I_n & 0 \end{array} \right), \quad (F_{kj}) = \left(\begin{array}{c|c} 0 & iI_n \\ \hline -iI_n & 0 \end{array} \right). \quad (1.2)$$

2. CARTAN-KIRICHENKO STRUCTURAL EQUATIONS OF AN AH -STRUCTURE

The form of the Levi-Civita connection ∇ is determined by the forms system $\{\omega_j^k\}$ on the space of the complex frames stratification over an almost Hermitian manifold. Similarly, the displacement form ω is determined by the forms system $\{\omega^k\}$. The Cartan structural equations of the stratification space over almost Hermitian manifold look as follows:

$$\begin{aligned} 1) \quad d\omega^k &= \omega_j^k \wedge \omega^j; \\ 2) \quad d\omega_j^k &= \omega_i^k \wedge \omega_j^i + \frac{1}{2} R_{jml}^k \omega^m \wedge \omega^l, \end{aligned} \quad (2.1)$$

where $\{R_{jml}^k\}$ are the components of the Riemannian curvature tensor (or of the Riemann-Christoffel tensor [13]). Here and further $k, j, m, l = 1, \dots, 2n$.

As J and g are the tensors of (1,1)- and (2,0)-type, respectively, and as $\nabla g = 0$, then the components of these tensors must satisfy the following system of differential equations:

$$\begin{aligned} 1) \quad dJ_j^k + J_l^k \omega_j^l - J_j^l \omega_l^k &= J_{j,l}^k \omega^l; \\ 2) \quad dg_{kj} + g_{lj} \omega_k^l + g_{kl} \omega_j^l &= 0, \end{aligned} \quad (2.2)$$

where $\{J_{j,l}^k\}$ are the components of ∇J . Taking into account (1.1) and (1.2), we can rewrite (2.2)₁ as follows:

$$1) \omega_b^a = -\frac{i}{2} J_{b,k}^a \omega^k; \quad 2) \widehat{\omega}_b^a = -\frac{i}{2} \widehat{J}_{b,k}^a \omega^k; \quad 3) J_{b,k}^a = 0; \quad 4) \widehat{J}_{b,k}^a = 0. \quad (2.3)$$

Similarly, from (2.2)₂ we obtain:

$$1) \omega_b^a + \omega_a^b = 0; \quad 2) \omega_b^a + \omega_a^{\widehat{b}} = 0; \quad 3) \omega_b^{\widehat{a}} + \omega_a^{\widehat{b}} = 0. \quad (2.4)$$

Here and further $a, b, c, d = 1, \dots, n$; $\widehat{a} = a + n$. Substituting (2.3) and (2.4) in Cartan structural equations (2.1), we get

$$\begin{aligned} d\omega^a &= \omega_b^a \wedge \omega^b + \omega_b^{\widehat{a}} \wedge \omega^{\widehat{b}} = \omega_b^a \wedge \omega^b - \frac{i}{2} J_{b,c}^a \omega^c \wedge \omega^{\widehat{b}} + \frac{i}{2} J_{[\widehat{b},c]}^a \omega^{\widehat{c}} \wedge \omega^{\widehat{b}}; \\ d\omega^{\widehat{a}} &= \omega_b^{\widehat{a}} \wedge \omega^b + \omega_b^{\widehat{a}} \wedge \omega^{\widehat{b}} = \omega_b^{\widehat{a}} \wedge \omega^{\widehat{b}} + \frac{i}{2} \widehat{J}_{b,c}^{\widehat{a}} \omega^{\widehat{c}} \wedge \omega^b - \frac{i}{2} \widehat{J}_{[b,c]}^{\widehat{a}} \omega^c \wedge \omega^b. \end{aligned} \quad (2.5)$$

We denote $\omega_k = g_{kj} \omega^j$. In particular, $\omega_a = \omega^{\widehat{a}} = \overline{\omega^a}$. Taking into account this fact as well as (2.4), we can rewrite (2.5) as follows:

$$\begin{aligned} d\omega^a &= \omega_b^a \wedge \omega^b + B^{ab}{}_c \omega^c \wedge \omega_b + B^{abc} \omega_b \wedge \omega_c; \\ d\omega_a &= -\omega_a^b \wedge \omega_b + B_{ab}{}^c \omega_c \wedge \omega^b + B_{abc} \omega^b \wedge \omega^c, \end{aligned} \quad (2.6)$$

where

$$B^{ab}{}_c = -\frac{i}{2} J_{b,c}^a; \quad B_{ab}{}^c = \frac{i}{2} \widehat{J}_{b,c}^{\widehat{a}}; \quad B^{abc} = \frac{i}{2} J_{[\widehat{b},c]}^a; \quad B_{abc} = -\frac{i}{2} \widehat{J}_{[b,c]}^{\widehat{a}}. \quad (2.7)$$

The functions $\{B^{ab}{}_c\}$, $\{B_{ab}{}^c\}$, $\{B^{abc}\}$, $\{B_{abc}\}$ serve as components of complex tensors of an almost Hermitian manifold (M^{2n}, J, g) [2], because, considering the differential continuations of (2.7), it is not difficult to see that

$$\begin{aligned} dB^{ab}{}_c + B^{ab}{}_d \omega_c^d - B^{db}{}_c \omega_d^a - B^{ad}{}_c \omega_d^b &= B^{ab}{}_{,k} \omega^k; \\ dB_{ab}{}^c - B_{ab}{}^d \omega_d^c + B_{db}{}^c \omega_a^d + B_{ad}{}^c \omega_b^d &= B_{ab}{}^c{}_{,k} \omega^k; \\ dB^{abc} - B^{dbc} \omega_d^a - B^{adc} \omega_d^b - B^{abd} \omega_d^c &= B^{abc}{}_{,k} \omega^k; \\ dB_{abc} + B_{dbc} \omega_a^d + B_{adc} \omega_b^d + B_{abd} \omega_c^d &= B_{abc}{}_{,k} \omega^k. \end{aligned}$$

Definition 2.1 ([4]). The tensors with the components $\{B^{ab}{}_c\}$ and $\{B_{ab}{}^c\}$ are called virtual Kirichenko tensors of first and second order, respectively.

Definition 2.2 ([4]). The tensors with the components $\{B^{abc}\}$ and $\{B_{abc}\}$ are called structural Kirichenko tensors of first and second order, respectively.

Definition 2.3. The equations (2.6) are called Cartan-Kirichenko structural equations of an almost Hermitian structure on the manifold M^{2n} .

We remark that according to (2.4)

$$J_{b,k}^a + J_{a,k}^b = 0.$$

So, we have

$$B^{ab}{}_c + B^{ba}{}_c = 0.$$

Similarly,

$$B_{ab}{}^c + B_{ba}{}^c = 0.$$

Thus, we have proved

Proposition 2.1. *The virtual Kirichenko tensors of an almost Hermitian manifold are skew-symmetric relative to the first pair of indices.*

From (2.7) we obtain the following result:

Proposition 2.2. *The structural Kirichenko tensors of an almost Hermitian manifold are skew-symmetric relative to the second pair of indices.*

Owing to the reality of ∇J , from the given definitions we have

Proposition 2.3. $B^{ab}{}_c = \overline{B_{ab}{}^c}$, $B^{abc} = \overline{B_{abc}}$.

3. THE MAIN RESULT

Let $\mathbf{O} \equiv \mathbf{R}^8$ be the Cayley algebra. As it is well-known [10], two non-isomorphic three-fold vector cross products are defined on it by means of the relations

$$P_1(X, Y, Z) = -X(\overline{Y}Z) + \langle X, Y \rangle Z + \langle Y, Z \rangle X - \langle Z, X \rangle Y,$$

$$P_2(X, Y, Z) = -(X\overline{Y})Z + \langle X, Y \rangle Z + \langle Y, Z \rangle X - \langle Z, X \rangle Y,$$

where $X, Y, Z \in \mathbf{O}$, $\langle \cdot, \cdot \rangle$ is the scalar product in \mathbf{O} , and $X \rightarrow \overline{X}$ is the conjugation operator. Moreover, any other three-fold vector cross product in the octave algebra is isomorphic to one of the two above-mentioned.

If $M^6 \subset \mathbf{O}$ is a six-dimensional oriented submanifold, then the induced almost Hermitian structure $\{J_\alpha, g = \langle \cdot, \cdot \rangle\}$ is determined by the relation

$$J_\alpha(X) = P_\alpha(X, e_1, e_2), \quad \alpha = 1, 2,$$

where $\{e_1, e_2\}$ is an arbitrary orthonormal basis of the normal space of M^6 at a point p , $X \in T_p(M^6)$ [10]. The point $p \in M^6$ is called general [12] if

$$e_0 \notin T_p(M^6) \not\subset L(e_0)^\perp,$$

where e_0 is the unit of Cayley algebra and $L(e_0)^\perp$ is its orthogonal supplement. A submanifold $M^6 \subset \mathbf{O}$, consisting only of general points, is called a general-type submanifold [12, 13].

Naturally, we can consider the following question:

Which of the sixteen Gray-Hervella classes of almost Hermitian structures can be represented on six-dimensional submanifolds of Cayley algebra?

A partial answer is known just for a six-dimensional special $M^6 \subset \mathbf{O}$, or the so-called Calabi submanifolds [13]. We recall that an almost Hermitian submanifold $M^6 \subset \mathbf{O}$ is called special [13] if

$$e_0 \notin T_p(M^6) \subset L(e_0)^\perp.$$

Such six-dimensional almost Hermitian submanifolds of the octave algebra were studied by A. Gray [8 – 10], E. Calabi [6], K. Yano and T. Sumitomo [17].

We answer this question in the case of general-type $M^6 \subset \mathbf{O}$, i.e. we shall prove the following

Main Theorem. *Just the eight Gray-Hervella classes of almost Hermitian structures can be represented on six-dimensional general-type submanifolds of Cayley algebra, namely:*

$$K, AK, NK, SH, QK, W_1 \oplus W_3, W_2 \oplus W_3, SK.$$

Proof. Let us use the characterization of Gray-Hervella classes of AH-structures in terms of Kirichenko tensors [2], as listed in Table 1.

Table 1
Characterization of Gray-Hervella classes of AH-structures
in terms of Kirichenko tensors [2]

Class	Condition
K	$B^{abc} = 0, B^{ab}_c = 0$
$NK = W_1$	$B^{abc} = -B^{bac}, B^{ab}_c = 0$
$AK = W_2$	$B^{(abc)} = 0, B^{ab}_c = 0$
$SH = W_3$	$B^{abc} = 0, B^{ab}_b = 0$
W_4	$B^{abc} = 0, B^{ab}_c = \alpha^{[a} \delta_c^{b]}$
$QK = W_1 \oplus W_2$	$B^{ab}_c = 0$
$H = W_3 \oplus W_4$	$B^{abc} = 0$
$VG = W_1 \oplus W_4$	$B^{abc} = -B^{bac}, B^{ab}_c = \alpha^{[a} \delta_c^{b]}$
$W_1 \oplus W_3$	$B^{abc} = -B^{bac}, B^{ab}_b = 0$
$W_2 \oplus W_3$	$B^{(abc)} = 0, B^{ab}_b = 0$
$W_2 \oplus W_4$	$B^{(abc)} = 0, B^{ab}_c = \alpha^{[a} \delta_c^{b]}$
$SK = W_1 \oplus W_2 \oplus W_3$	$B^{ab}_b = 0$
$G1 = W_1 \oplus W_3 \oplus W_4$	$B^{abc} = -B^{bac}$
$G2 = W_2 \oplus W_3 \oplus W_4$	$B^{(abc)} = 0$
$W_1 \oplus W_2 \oplus W_4$	$B^{ab}_c = \alpha^{[a} \delta_c^{b]}$
W	no condition

Now, we write out the Cartan-Kirichenko structural equations for six-dimensional general-type almost Hermitian submanifolds of Cayley algebra [13]:

$$\begin{aligned} d\omega^a &= \omega_b^a \wedge \omega^b + \frac{1}{\sqrt{2}}\varepsilon^{ah[b}D_{h^c]}\omega_b \wedge \omega_c + \frac{1}{\sqrt{2}}\varepsilon^{abh}D_{hc}\omega^c \wedge \omega_b; \\ d\omega_a &= -\omega_a^b \wedge \omega_b + \frac{1}{\sqrt{2}}\varepsilon_{ah[b}D^h_{c]}\omega_b \wedge \omega_c + \frac{1}{\sqrt{2}}\varepsilon_{abh}D^{hc}\omega_c \wedge \omega^b. \end{aligned} \quad (3.1)$$

Here $\varepsilon^{abc} = \varepsilon_{123}^{abc}$, $\varepsilon_{abc} = \varepsilon_{abc}^{123}$ are the components of Kronecker tensor of third order [16]:

$$\begin{aligned} D_{cj} &= \mp T_{cj}^8 + iT_{cj}^7, & D_{\widehat{c}j} &= \mp T_{\widehat{c}j}^8 - iT_{\widehat{c}j}^7; \\ D_h^c &= D_{h\widehat{c}}, & D^h_c &= D_{\widehat{h}c}, & D^{hc} &= D_{\widehat{h}c}, \end{aligned} \quad (3.2)$$

where $\{T_{kj}^\varphi\}$, $\varphi = 7, 8$, are components of the configuration tensor (using Gray's terminology [8] or of the Euler curvature tensor [7]). Now, we assume that the indices a, b, c, d, h range from 1 to 3; the indices k, j range from 1 to 6, and we set $\widehat{a} = a + 3$.

Comparing (2.6) and (3.1), we get the following relations for virtual Kirichenko tensors of an almost Hermitian $M^6 \subset \mathbf{O}$ [5]:

$$B_{ab}{}^c = \frac{1}{\sqrt{2}}\varepsilon_{abh}D^{hc}, \quad B^{ab}{}_c = \frac{1}{\sqrt{2}}\varepsilon^{abh}D_{hc}.$$

From (3.2) it follows that the tensor D_{hb} is symmetric relative to the indices h and b ; ε^{abh} is skew-symmetric relative to these indices. Hence, $\varepsilon^{abh}D_{hb} = 0$, and therefore

$$B^{ab}{}_b = \frac{1}{\sqrt{2}}\varepsilon^{abh}D_{hb} = 0.$$

So, for an arbitrary almost Hermitian structure induced by means of a three fold vector cross product on six-dimensional submanifolds of Cayley algebra, the following identity is fulfilled:

$$B^{ab}{}_b \equiv 0.$$

Consequently, as it is clear from the given table, all these almost Hermitian structures must be semi-Kählerian (*SK*). The class of *SK*-structures contains only eight classes of *AH*-structures, namely:

$$K, NK, AK, SH, QK, W_1 \oplus W_3, W_2 \oplus W_3, SK. \quad \square$$

We remark that this result is similar to A. Gray's conclusion that a six-dimensional special submanifold $M^6 \subset \mathbf{O}$ is semi-Kählerian [9, 11]. So, all six-dimensional almost Hermitian submanifolds of Cayley algebra are semi-Kählerian, i.e. they belong to one of the eight above-mentioned Gray-Hervella classes of *AH*-manifolds.

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Smolensk University of Humanities
 2, Gertsen str., 214014 Smolensk
 RUSSIA
 E-mail: banaru@keytown.com